

# Stress measurements in ZrB<sub>2</sub>–SiC composites using Raman spectroscopy and neutron diffraction

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## Abstract

Raman spectroscopy and neutron diffraction were used to study the stresses generated in zirconium diboride–silicon carbide (ZrB<sub>2</sub>–SiC) ceramics. Dense, hot pressed samples were prepared from ZrB<sub>2</sub> containing 30 vol%  $\alpha$ -SiC particles. Raman patterns were acquired from the dispersed SiC particulate phase within the composite and stress values were calculated to be 810 MPa. Neutron diffraction patterns were acquired for the ZrB<sub>2</sub>–SiC composite, as well as pure ZrB<sub>2</sub> and SiC powders during cooling from  $\sim$ 1800 °C to room temperature. A residual stress of 775 MPa was calculated as a function of temperature by comparing the lattice parameter values for ZrB<sub>2</sub> and SiC within the composite to those of the individual powders. The temperature at which stresses began to accumulate on cooling was found to be  $\sim$ 1400 °C based on observing the deviation in lattice parameters between pure powder samples and those of the composite.

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## 1. Introduction

Zirconium diboride (ZrB<sub>2</sub>) belongs to a class of materials known as ultra-high temperature ceramics, which exhibit melting temperatures in excess of 3000 °C.<sup>1,2</sup> The refractory diborides are generally resistant to chemical attack; this in addition to its high melting point has led to interest in ZrB<sub>2</sub> as a material for molten metal crucibles.<sup>3</sup> Along with high melting temperature and chemical resistivity, zirconium diboride composites, specifically those produced using silicon carbide (SiC) additives, display a wide variety of desirable properties. ZrB<sub>2</sub>–SiC composites have been shown to have strengths in excess of 1000 MPa<sup>4–6</sup> and hardness values of 22 GPa and higher.<sup>4,5,7</sup> Zirconium diboride's unusual combination of properties have resulted in interest for uses in refractory linings,<sup>8–10</sup> microelectronics,<sup>11</sup> electrodes,<sup>12–14</sup> cutting tools,<sup>3,15</sup> as well as for use in future aerospace vehicles.<sup>16–18</sup>

The addition of SiC to ZrB<sub>2</sub> has been shown to improve its strength,<sup>4–6,19–21</sup> oxidation resistance,<sup>21–23</sup> and fracture toughness<sup>4,19</sup>; all important factors for potential use in aerospace

applications. These effects are believed to be due in part to the stresses generated in and around SiC particles as these materials are cooled from the processing temperature. The room temperature coefficient of thermal expansion (CTE) of ZrB<sub>2</sub> is approximately 5.2 ppm/K<sup>2,24</sup> while the CTE of the 6H polytype of  $\alpha$ -SiC is approximately 3.3 ppm/K.<sup>24</sup> As a result, as these materials are cooled from their final densification temperature (typically >1800 °C), compressive stresses are generated in the SiC particles and corresponding tensile stresses are generated in the ZrB<sub>2</sub> matrix as it tries to shrink at a faster rate than the SiC. Raman spectroscopy is one method for measuring stresses in solids. Unfortunately ZrB<sub>2</sub> is not Raman active; however, SiC is. The Raman patterns obtained from Raman active materials provide information similar to other spectroscopic techniques in that the pattern is unique to that specific material. The locations of the peaks present in a Raman pattern are also sensitive to stress.<sup>25–29</sup> While no general relationship equating peak shift to stress has been established, calibration curves have been experimentally obtained for a variety of materials making it possible to correlate peak shift to the stress that material is experiencing.<sup>28,30,31</sup> One limitation of Raman spectroscopy is that it is surface sensitive for opaque materials such as SiC. Neutron diffraction, on the other hand, is capable of providing lattice spacing data from the bulk of a sample, not just at the surface. As a result, neutron

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diffraction can be used to measure the internal stresses within a material by comparing the measured peak positions to the known peak positions from stress free materials.

The goal of the current study was to measure the residual stresses that develop in hot pressed ZrB<sub>2</sub>–SiC composites during cooling from the final densification temperature. This was accomplished using both Raman spectroscopy and neutron diffraction. Raman data was collected on individual SiC particles (~2 μm in size) within the composite. Special samples had to be prepared in order to collect neutron diffraction data. While naturally occurring boron consists of 80% <sup>11</sup>B and 20% <sup>10</sup>B the high neutron absorption cross section of <sup>10</sup>B (3835 b)<sup>32</sup> makes neutron diffraction impractical. The absorption cross section of <sup>11</sup>B however is only 0.0055 b.<sup>32</sup> In light of this, samples were produced via reactive hot pressing using ZrH<sub>2</sub> and <sup>11</sup>B powder precursors to produce ZrB<sub>2</sub> composites suitable for neutron diffraction studies. Three samples were then prepared for neutron diffraction: pure Zr<sup>11</sup>B<sub>2</sub>, pure SiC, and a composite consisting of 70 vol% Zr<sup>11</sup>B<sub>2</sub> with 30 vol% SiC. Neutron diffraction patterns were gathered on all three samples at temperature intervals ranging from room temperature to 1750 °C to determine the total magnitude of stress present as well as to determine the temperature at which stresses begin to accumulate during cooling.

## 2. Experimental procedure

### 2.1. Materials and processing

Two powder preparation methods were used in this study. The samples prepared for Raman characterization were produced from ZrB<sub>2</sub> powder (H.C. Starck, Grade B, Goslar, Germany) and SiC powder (H.C. Starck, Grade UF-25). The UF-25 SiC is an alpha-SiC consisting primarily of the 6H polytype. A mixture of 70 vol% ZrB<sub>2</sub> and 30 vol% SiC powder was then attrition milled using WC milling media in a Teflon lined vessel and acetone for a duration of 2 h. Following attrition milling the acetone was removed from the milled slurry using rotary evaporation. The dried powder was then passed through a 60-mesh sieve in order to facilitate uniform die filling later in the process. The samples for neutron diffraction were prepared reactively using ZrH<sub>2</sub> (Alfa Aesar, Grade Z-1038, Ward Hill, MA, USA), boron (Eagle Picher Boron Products, Isotopic enriched boron 11 metal, Quapaw, OK, USA) and the same SiC as above. These powders were ball milled in a polyethylene jar for 12 h using acetone and WC milling media in place of attrition milling for 2 h as was done for the other samples in an effort to affect the same particle size reduction while minimizing WC contamination. The solvent was then removed using rotary evaporation and the dried powders were passed through a 60-mesh sieve.

Both compositions were hot pressed in 44 mm diameter circular graphite dies. The hot press dies were lined with graphoil and the ends of the rams, which would be in contact with the powder, were coated with boron nitride spray (Cerac, SP-108, Milwaukee, WI, USA). After the powder was loaded into the die, it was cold pressed to ~6 MPa. All densification in this study was performed using a resistively heated graphite element hot press

(Thermal Technology Inc., Model HP20-3060-20, Santa Rosa, CA, USA). Below 1650 °C, samples were heated under mild vacuum (~20 Pa). The sample being prepared reactively was heated at ~5 °C/min up to 700 °C for the ZrH<sub>2</sub> and <sup>11</sup>B to react to form ZrB<sub>2</sub>. The non-reactive sample was heated directly to the first hold temperature (1450 °C) at a rate of ~75 °C/min. For both compositions, the powder compact was then heated successively to hold temperatures of 1450 °C and 1650 °C to facilitate the removal of oxide species from the surfaces of the powders. After the 1650 °C hold the furnace atmosphere was changed from vacuum to argon, a hydraulic ram pressure of 32 MPa was applied, and the temperature was raised at ~60 °C/min to 1900 °C. The movement of the hot press rams was monitored and used to determine when densification had ceased. Once ram travel had ceased, the furnace power was turned off and the furnace was allowed to cool (~50 °C/min from 1900 °C to below 600 °C). Load on the sample was maintained until the temperature fell to 1650 °C and was then removed. Powder specimens for neutron diffraction were prepared by reacting ZrH<sub>2</sub> and <sup>11</sup>B powders up to 1450 °C and crushing the resulting material to –325 mesh for the Zr<sup>11</sup>B<sub>2</sub> while the SiC was used as received.

After hot pressing, billets were removed from the die and surface ground in order to make the surfaces flat and parallel as well as remove any boron nitride adhering to the specimen surfaces. Each billet was then cut into bars ~2.5 mm in width. Finally the bars were ground to obtain final cross sectional dimensions of 1.5 mm × 2 mm with a 400-grit finish. All cutting and grinding was done with diamond abrasive wheels on an automated surface grinder (Chevalier, FSG-3A818, Santa Fe Springs, CA). Specimens prepared for Raman measurements and microscopy were polished to a 0.25 μm diamond finish.

### 2.2. Characterization

Archimedes' density measurements were made for each billet. After the dry weight was recorded, samples were saturated by submersion in distilled water and placing them under vacuum for 12 h. The saturated and suspended weights were then measured to calculate a final density. Microstructural analysis was performed using a Hitachi S570 scanning electron microscope (Tokyo, Japan). Particle size was determined using ImageJ software by analyzing ~500 grains. X-ray diffraction analysis was performed using a Philips X-Pert Pro diffractometer (Westborough, MA, USA).

Stress values were obtained using both Raman spectroscopy and neutron diffraction. Raman measurements were made using a Horiba LabRAM ARAMIS spectrometer (Edison, NJ, USA) using a 633 nm HeNe laser and a 1 μm spot size. Before the dataset was obtained the instrument was calibrated using a silicon standard, and data was collected up to 1000 wave numbers. Raman patterns were gathered only from the SiC within the ZrB<sub>2</sub>–SiC composite material due to ZrB<sub>2</sub> being Raman inactive. Data was gathered from 25 separate SiC particulates within the microstructure. Neutron diffraction was performed using the time of flight (TOF) method in the Lujan Center at Los Alamos National Lab using the Spectrometer for Materials Research at Temperature and Stress (SMARTS). Specimens for neutron

Table 1

Final density and SiC particle size of the reactively and traditionally produced specimens.

Composition	Theoretical density (g/cm <sup>3</sup> )	Archimedes density (g/cm <sup>3</sup> )	Average SiC particle size (μm)
30 vol% SiC 70 vol% ZrB <sub>2</sub>	5.24	5.32	3.3
30 vol% SiC 70 vol% Zr <sup>11</sup> B <sub>2</sub>	5.25	5.38	2.0

diffraction consisted of 5 bars (1.5 mm × 2 mm × 15 mm) contained in an alumina crucible; the crucible was placed inside a tungsten element furnace situated in the path of the neutron beam. The specimens were heated to a maximum temperature of 1750 °C under vacuum (1 × 10<sup>-4</sup> Torr) at a rate of 20 °C/min. Neutron diffraction patterns were collected at various temperatures upon both heating and cooling. Temperature was allowed to equilibrate for 15 min prior to each measurement. Data was then gathered for approximately 1 h at each temperature. Diffraction data was collected for a composite material consisting of 70 vol% Zr<sup>11</sup>B<sub>2</sub> and 30 vol% SiC as well as for the individual Zr<sup>11</sup>B<sub>2</sub> and SiC powders.

### 3. Results and discussion

Final densities and SiC particle sizes were measured for both the ZrB<sub>2</sub>-SiC and Zr<sup>11</sup>B<sub>2</sub>-SiC composite materials (Table 1) to ensure that microstructures and densities were similar. Both materials were hot pressed to full density as measured by Archimedes' technique. As can be seen in Table 1 both materials actually reached a slightly higher density than predicted based on the nominal composition of 70 vol% ZrB<sub>2</sub> and 30 vol% SiC. This was due to the incorporation of small amounts of WC milling media used; approximately 2 wt% WC was in the ZrB<sub>2</sub>-SiC composite and 4 wt% was in the Zr<sup>11</sup>B<sub>2</sub>-SiC composite as determined by the mass loss of the milling media. The microstructures of both materials consisted of a ZrB<sub>2</sub> matrix containing a fine, well-dispersed SiC second phase (Fig. 1). The SiC particulate size in the reaction processed specimen (~2 μm) was slightly smaller than for the conventionally hot pressed specimen (~3 μm). This is consistent with previously reported results.<sup>7</sup> Given the similarity in microstructures there should not be a significant difference in stress evolution during processing between the two materials.

A Raman pattern (Fig. 2) obtained from the SiC particles in the hot pressed ZrB<sub>2</sub>-30 vol% SiC specimen shows measurable peak shifts in the wave number for both the transverse optic (TO) and linear optic (LO) peaks. The TO and LO peaks have been shifted by 2.5 cm<sup>-1</sup> and 3.9 cm<sup>-1</sup>, respectively, to higher wave numbers. Fig. 2 illustrates how the peaks in the Raman pattern of stressed SiC shift away from the standard, unstressed locations. This phenomenon has been studied for a number of materials and can be used to calculate the internal stress of the material being analyzed.<sup>26,31,33</sup>

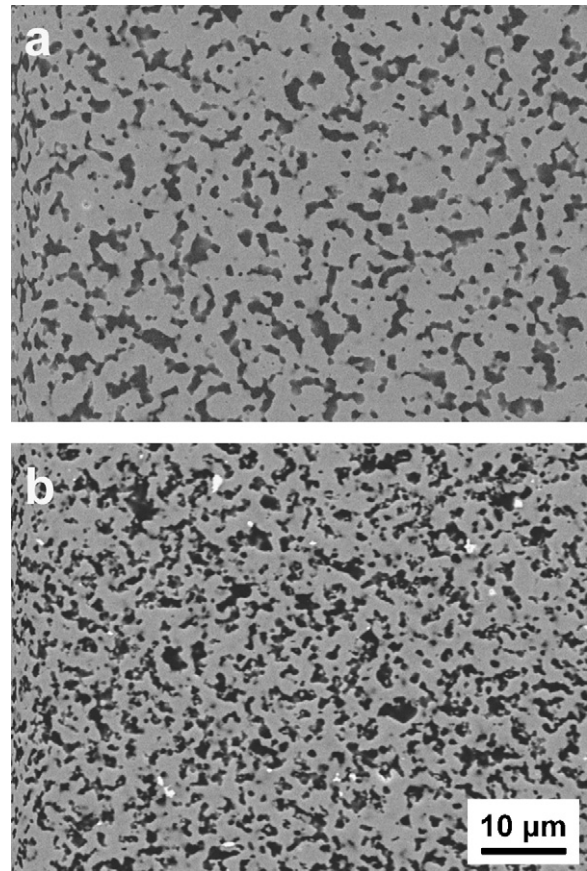


Fig. 1. Microstructural comparison of conventionally hot pressed ZrB<sub>2</sub>-SiC (a) with reaction hot pressed Zr<sup>11</sup>B<sub>2</sub>-SiC (b).

Liu and Vohra<sup>26</sup> analyzed Raman peak shifting of 6H SiC as a function of applied hydrostatic pressure from ambient up to 90 GPa compressive stress. Further, they were able to develop an equation (Eq. (1)) to correlate the magnitude of peak shift to the internal stress in 6H SiC.<sup>26</sup>

$$\omega_{\text{TO}}(\text{cm}^{-1}) = 789.2 + 3.11P - 0.009P^2 \quad (1)$$

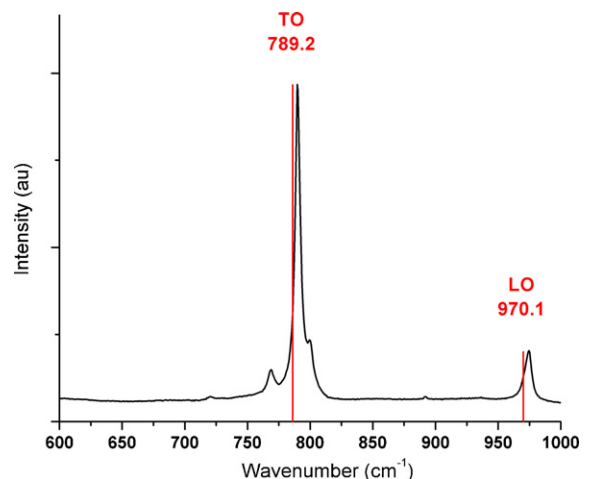


Fig. 2. Raman pattern for 6H SiC from a hot pressed ZrB<sub>2</sub>-SiC composite with unstressed peak positions identified with lines.<sup>26</sup>

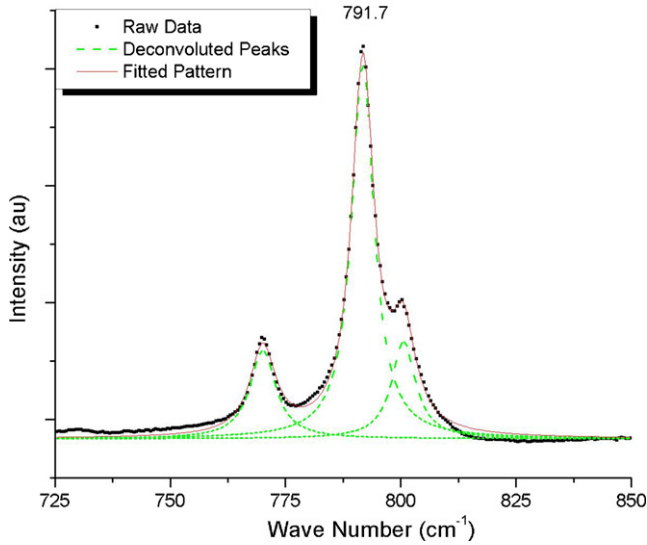


Fig. 3. Representative raw Raman spectrum from an individual SiC particle as well as deconvoluted peak fits to determine true peak position.

In this equation  $\omega_{TO}$  refers to the peak position of the transverse optical (TO) peak, and  $P$  is the internal stress in GPa. Using the unstressed position of the TO peak ( $789.2 \text{ cm}^{-1}$ ),<sup>26</sup> and combining that term with the  $\omega_{TO}$  term and solving for  $P$ , a second equation which allows for the calculation of internal stress by inputting the amount of peak shift ( $\Delta TO$ ) is obtained (Eq. (2)).

$$P = \frac{3.11 - \sqrt{9.6721 - 0.036(\Delta TO)}}{0.018} \quad (2)$$

A representative Raman pattern (Fig. 3) acquired from an individual SiC particle in the conventionally hot pressed  $\text{ZrB}_2$ -30 vol% SiC material shows a TO peak shift of  $2.5 \text{ cm}^{-1}$ . The TO peak was focused on for this study due to the strength of the peak and the ability to obtain accurate peak fits. When two or more peaks overlap one another, the resultant peak positions are affected due to the additive effect of the multiple peaks. Therefore, it is necessary to deconvolute the peaks to obtain the true peak positions. Fig. 3 is a plot of raw Raman data representative of the dataset along with deconvoluted peak fits to determine the true center of the TO peak.

By inputting the TO peak shift ( $2.5 \text{ cm}^{-1}$ ) for  $\Delta TO$  in Eq. (2), the internal stress in a representative SiC particle was calculated to be approximately 810 MPa (compressive). Using that stress value ( $\sigma$ ), along with the elastic moduli ( $E$ ) and Poisson's ratios ( $\nu$ ) of SiC and  $\text{ZrB}_2$ , the difference in thermal expansion coefficient ( $\Delta\alpha$ ) between SiC and  $\text{ZrB}_2$ , and the volume fraction of particulate phase ( $V_p$ ), it is possible to calculate the temperature at which stresses began to accumulate upon cooling (Eq. (3)). In this equation the subscripts p and m designate particle and matrix respectively.<sup>34</sup>

$$\Delta T = \frac{\sigma(0.5(1+\nu_m) + (1-2\nu_m)V_p/E_m(1-V_p) + 1-2\nu_p/E_p)}{\Delta\alpha} \quad (3)$$

While the CTE of both SiC and  $\text{ZrB}_2$  change as a function of temperature, the difference in CTE is  $\sim 2 \text{ ppm/K}$  over the entire processing temperature range. By inputting that value for  $\Delta\alpha$ ,  $415 \text{ GPa}^{35}$  for  $E_p$ ,  $489 \text{ GPa}^2$  for  $E_m$ , 0.14 and 0.16 for  $\nu_m^{36}$  and  $\nu_p^{35}$  respectively, and the measured compressive stress of 810 MPa for  $\sigma$ ,  $\Delta T$  is found to be  $1595 \text{ }^\circ\text{C}$ . Taking into account that the stress was measured at room temperature, the measurement suggests that stresses started to accumulate at approximately  $1620 \text{ }^\circ\text{C}$  during cooling. It is also important to note here that Raman spectroscopy is a surface sensitive characterization technique. Therefore, the stresses measured on a free surface of the sample could be different than those that exist in the bulk of the material.

Neutron diffraction, on the other hand, allows in-situ measurement of the internal strains in the bulk of the material as the specimen is heated to near the processing temperature and then cooled. Fig. 4 illustrates the lattice parameters for the (006) SiC plane between a powder specimen (i.e., no residual stress due to CTE mismatch with a matrix phase) and SiC in the  $\text{Zr}^{11}\text{B}_2$ -30 vol% SiC specimen. Diffraction data was collected as the specimen was cooled from  $1750 \text{ }^\circ\text{C}$  to room temperature. The CTE was measured for the SiC powder based on the shift in lattice parameter as a function of temperature. Fitting the data with a 3rd order polynomial shows good agreement with data in the literature<sup>24,35</sup> and results in a room temperature CTE of  $3.7 \text{ ppm}/^\circ\text{C}$ .

At temperatures of  $1500 \text{ }^\circ\text{C}$  and higher, the lattice parameters of the powder and the composite overlapped one another as is expected when both materials are stress free. However, at some temperature between  $1500 \text{ }^\circ\text{C}$  and  $1200 \text{ }^\circ\text{C}$  the lattice parameters of the two materials begin to deviate. The major reason for the two lattice parameters to differ from each other as a function of temperature was the stress generated in the composite specimen during cooling due to its CTE mismatch with the  $\text{ZrB}_2$  matrix. Knowing that at  $1200 \text{ }^\circ\text{C}$  and below the lattice parameters differed between powder and composite, the

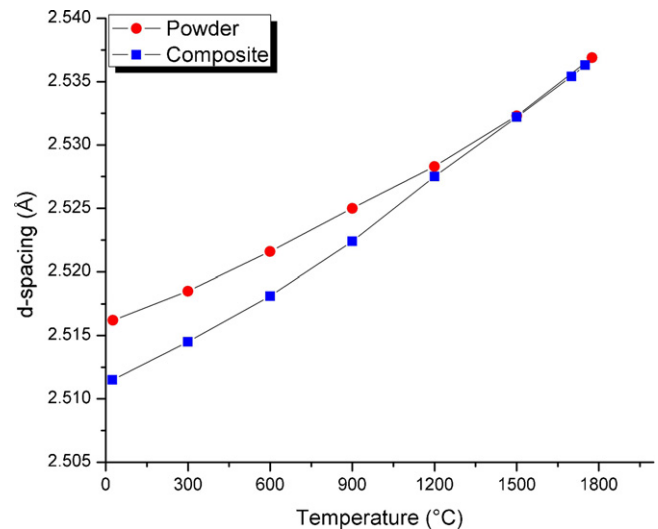


Fig. 4. The (006) lattice parameter of both SiC powder and SiC within the  $\text{Zr}^{11}\text{B}_2$ -SiC composite as a function of temperature during cooling from  $1775 \text{ }^\circ\text{C}$ .

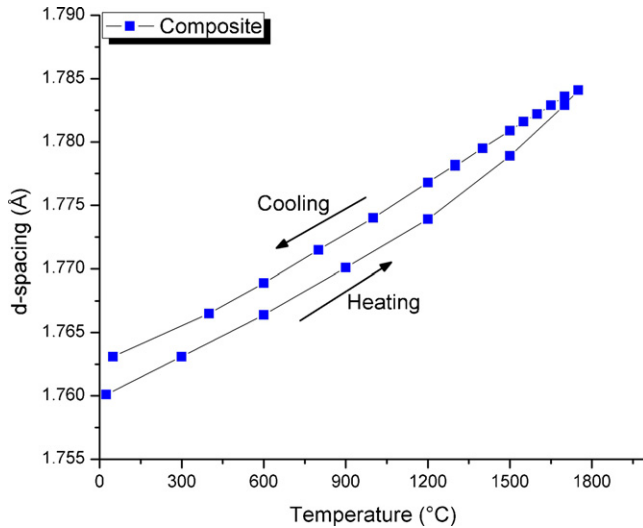


Fig. 5. The (002) lattice parameter for  $Zr^{11}B_2$  in the composite deviating from itself between heating and cooling.

composite lattice parameter data below  $1200^\circ C$  was fit with a second order polynomial and extrapolated to higher temperature. This resulted in the extrapolated fit intercepting the powder lattice parameter at  $\sim 1400^\circ C$ . The room temperature lattice parameter difference between the powder and composite specimens was measured to be  $0.0047 \text{ \AA}$ , which equates to a strain of  $0.187\%$ . By relating the strain to stress using the elastic modulus, the stress within the SiC was found to be  $\sim 775 \text{ MPa}$  (compressive).

In order to determine if the  $\Delta T$  values calculated from the Raman stress measurements were reasonable, a similar calculation was performed using the stress of  $775 \text{ MPa}$  calculated from neutron diffraction and the  $\Delta T$  range known from the lattice parameter measurements. Inputting  $775 \text{ MPa}$  into Eq. (3) yields a temperature where stresses begin to accumulate of  $\sim 1550^\circ C$ , which is higher than the  $1400^\circ C$  predicted by neutron diffraction. However, further investigation into the neutron diffraction data revealed that more than just a CTE mismatch between  $ZrB_2$  and SiC affected the lattice parameter shifts within the composite.

The lattice parameters for  $ZrB_2$  in the composite should be shifted as compared to those of the standard powder due to internal stresses. However, they should be self-consistent between heating and cooling cycles. The (002) lattice parameter measured for the  $ZrB_2$  within the composite (Fig. 5), differed at any given temperature during heating and cooling. The lattice parameter became larger as temperature increased, as expected due to thermal expansion. However during cooling, the lattice parameter decreased, but remained larger than it was on heating over the entire temperature range. Fig. 5 illustrates how the (002) lattice parameter of  $ZrB_2$  within the composite changed during heating and cooling. The deviation between heating and cooling likely indicated a compositional change within the composite that was shifting the lattice parameters after prolonged exposure to high temperatures.

Microscopy analysis was used to investigate the compositional changes in the  $Zr^{11}B_2$ -30 vol% SiC specimen after

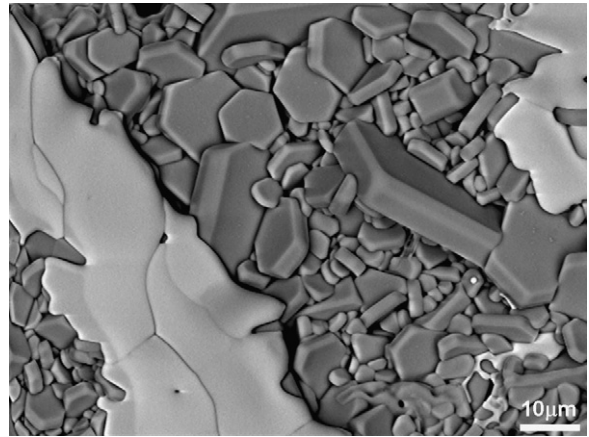


Fig. 6. Backscatter SEM image of the surface of a sample exhibiting a high z-contrast second phase generated during the neutron diffraction experiment.

elevated temperature neutron diffraction. The SEM images revealed the presence of a third phase (Fig. 6) in addition to  $ZrB_2$  and SiC. This third phase presumably evolved during the high temperature portion of the experiment. Energy dispersive spectroscopy (EDS) indicated that the third phase was rich in tungsten and silicon. Fig. 6 is a backscattered SEM image of the surface of the specimen showing the presence of a high z-contrast phase with respect to  $ZrB_2$ . X-ray diffraction performed on that same specimen also indicated the presence of the third phase, which exhibited diffraction peaks consistent with  $WSi_2$  (Fig. 7). Fig. 7 illustrates the presence of three peaks consistent with  $WSi_2$  in what is otherwise a typical  $ZrB_2$  diffraction pattern. Neither XRD nor SEM indicated the presence of this phase prior to the neutron diffraction experiment. Based on the introduction of WC into the specimen during milling, it appears that the W went into solid solution during processing, and stayed in solution after the relatively rapid cooling ( $\sim 100^\circ C/min$ ) following densification. During the neutron diffraction experiment however, the specimen was held at various temperatures below

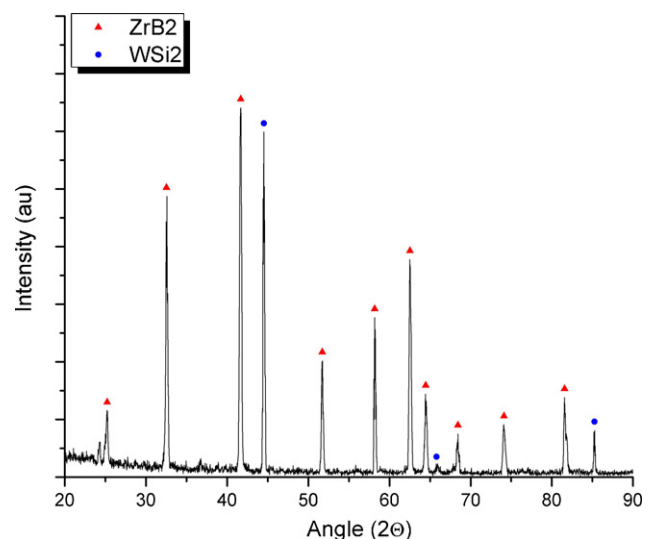


Fig. 7. X-ray diffraction spectrum from the surface of a sample exhibiting a second phase consistent with  $WSi_2$ .

the densification temperature for over an hour at each temperature while the neutron diffraction pattern was being collected. It appears that the WC (or just the W) came out of solution and reacted with the SiC to form  $\text{WSi}_2$ . The exsolution of WC (or W) from  $\text{ZrB}_2$  would be expected to increase the lattice parameters as measured via neutron diffraction since W has a smaller covalent radius than Zr.<sup>37</sup>

The stress calculated from Raman spectroscopy peak shifts (810 MPa) was slightly higher than that calculated from neutron diffraction analysis (775 MPa). This is counterintuitive based on the notion that the particles analyzed using Raman were on a free surface, which would presumably reduce stress compared to the bulk by releasing part of the constraint on the particle. While neutron diffraction has been extensively used to measure the internal stresses of composite materials,<sup>38–41</sup> it has also been shown that the measured lattice parameter shifts expected due to CTE mismatch can be altered by changes in chemistry.<sup>42</sup> If the third phase had the same effect on the SiC as on the  $\text{ZrB}_2$ , that of increasing the lattice parameter, it would mask some of the contraction in lattice parameter due to compressive stress and could result in a lower measured stress. In light of this, further neutron diffraction experiments could benefit from production of samples without any WC impurity to confirm the SiC lattice parameter data and collect complementary data for lattice parameter changes and residual stresses in  $\text{ZrB}_2$ .

#### 4. Conclusions

Raman spectroscopy and neutron diffraction were used to determine the stresses generated in SiC particles within  $\text{ZrB}_2$ -SiC composites during processing. Peak shifts measured by Raman spectroscopy were consistent with a compressive stress of  $\sim 810$  MPa at room temperature, which corresponds to a zero stress temperature of  $\sim 1620$  °C during cooling from the densification temperature.  $\text{Zr}^{11}\text{B}_2$ -SiC composites were successfully fabricated using  $^{11}\text{B}$  isotope. Lattice parameters for the composite were measured using neutron diffraction from room temperature up to 1750 °C. Lattice parameter changes in the SiC indicated a compressive stress of  $\sim 775$  MPa and a zero stress temperature of  $\sim 1400$  °C. The discrepancy between  $\text{Zr}^{11}\text{B}_2$  lattice parameters on heating and cooling, as well as subsequent XRD and SEM analysis, suggest that the WC impurity introduced during milling caused compositional changes and thus altered lattice parameters during the experiment.

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